

# Shape and topology optimisation subject to 3D nonlinear magnetostatics - part 1: sensitivity analysis

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# Outline

- Model Problem
- Topological Derivative
- Averaged adjoint formalism
- Application to model problem

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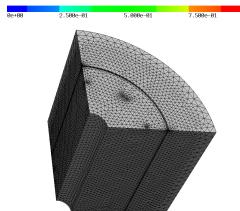
## Quasilinear Model Problem

$$\min_{\Omega \in \mathcal{A}(D)} J(\Omega) = \int_{\Gamma_0} |\mathbf{curl} \mathbf{u} \cdot \mathbf{n} - B_d^n|^2 dS_x$$

$$\text{s.t. } \mathbf{u} \in \mathbf{V} : \int_D \nu_\Omega(x, |\mathbf{curl} \mathbf{u}|) \mathbf{curl} \mathbf{u} \cdot \mathbf{curl} \mathbf{v} dx = \langle \mathbf{F}, \mathbf{v} \rangle \quad \text{for all } \mathbf{v} \in \mathbf{V}.$$

Here,

- $\nu_\Omega(x, s) = \chi_\Omega(x)\nu_2(s) + \chi_{D \setminus \bar{\Omega}}(x)\nu_1(s)$ ,
- $\langle \mathbf{F}, \mathbf{v} \rangle := \int_{\Omega_c} \mathbf{J} \cdot \mathbf{v} dx + \int_{\Omega_m} \nu_m \mathbf{M} \cdot \mathbf{curl} \mathbf{v} dx \quad \text{for } \mathbf{v} \in \mathbf{V}$ ,
- $\mathbf{V} := H_0(D, \mathbf{curl}) \cap H(D, \mathbf{div} = 0)$ ,
- $\nu_1, \nu_2 : \mathbf{R}_0^+ \rightarrow \mathbf{R}^+$  satisfy
  - $s \mapsto \nu_i(s)s, i = 1, 2$  is Lipschitz continuous and strongly monotone
  - $\nu_i \in C^2(\mathbf{R}_0^+)$ ,  $\nu_i'(0) = 0$ , and that there is a constant  $c$  such that for all  $s \in \mathbf{R}_0^+$ ,  $\nu_i'(s) \leq c$  and  $\nu_i''(s) \leq c$



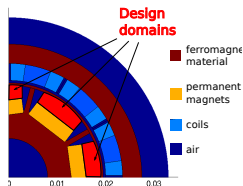
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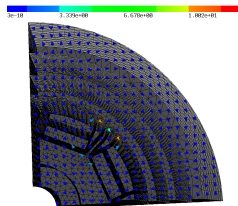
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# Topological Derivative: Definition

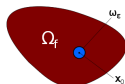
## Idea:

Sensitivity of  $\mathcal{J} = \mathcal{J}(\Omega) = J(\Omega, u(\Omega))$  w.r.t. insertion of hole  $\omega_\varepsilon = x_0 + \varepsilon\omega$  ( $\omega$  e.g. dots unit disk)

## Definition (Topological derivative)

Let  $D \subset \mathbf{R}^3$  be an open set and  $\Omega \subset D$  an open subset. Let  $\omega \subset \mathbf{R}^3$  be open with  $0 \in \omega$ . Define for  $z \in \mathbf{R}^3$ ,  $\omega_\varepsilon(z) := z + \varepsilon\omega$ . Then the topological derivative of  $J$  at  $\Omega$  at the point  $z \in \mathbf{R}^3$  is defined by

$$dJ(\Omega)(z) = \begin{cases} \lim_{\varepsilon \searrow 0} \frac{J(\Omega \setminus \omega_\varepsilon(z)) - J(\Omega)}{|\omega_\varepsilon(z)|} & \text{if } z \in \Omega, \\ \lim_{\varepsilon \searrow 0} \frac{J(\Omega \cup \omega_\varepsilon(z)) - J(\Omega)}{|\omega_\varepsilon(z)|} & \text{if } z \in D \setminus \bar{\Omega}. \end{cases}$$



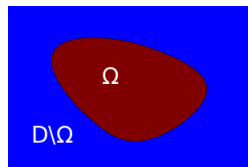
$G^{f \rightarrow \text{air}}$ : air inside ferromagnetic material



$G^{\text{air} \rightarrow f}$ : ferromagnetic material inside air



## Topological Derivative: Example

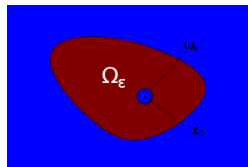


**Example:**  $\mathcal{J}(\Omega) := |\Omega|$ ,  $\omega_\varepsilon = x_0 + \varepsilon\omega$ ,  $\omega \in \mathbf{R}^d$ ,  $0 \in \omega$ ,

### Remark

Topological derivative can be seen as a semidifferential on the space of characteristic functions.

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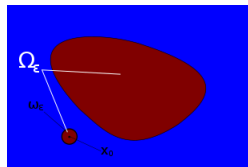
1  $x_0 \in \Omega$ ,  $\Omega_\varepsilon = \Omega \setminus \overline{\omega_\varepsilon}$  :

$$\frac{1}{|\omega_\varepsilon|} (\mathcal{J}(\Omega_\varepsilon) - \mathcal{J}(\Omega)) = -\frac{1}{|\omega_\varepsilon|} \int_{\omega_\varepsilon} f \, dx \rightarrow -f(x_0),$$

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2  $x_0 \in D \setminus \overline{\Omega}$ ,  $\Omega_\varepsilon = \Omega \cup \omega_\varepsilon$  :

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### Remark

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## Definition

Let  $X$  and  $Y$  be vector spaces and  $\tau > 0$ . A parametrised Lagrangian (or short Lagrangian) is a function

$$(\varepsilon, u, q) \mapsto L(\varepsilon, u, q) : [0, \tau] \times X \times Y \rightarrow \mathbf{R},$$

satisfying for all  $(\varepsilon, u) \in [0, \tau] \times X$ ,

$$q \mapsto L(\varepsilon, u, q) \quad \text{is affine on } Y.$$

## Example

Let  $X = Y = H_0^1(\Omega)$  and

$$L(\varepsilon, u, q) := \int_{\Omega} u^2 \, dx + \int_{\Omega} \nabla u \cdot \nabla q - \varepsilon f q \, dx.$$

## State and averaged adjoint state

### Definition (Perturbed states)

For  $\varepsilon \in [0, \tau]$  we define the (perturbed) state equation by: find  $u_\varepsilon \in X$ , such that

$$\partial_q L(\varepsilon, u_\varepsilon, 0)(\varphi) = 0 \quad \text{for all } \varphi \in Y.$$

The set of perturbed states is denoted  $E(\varepsilon)$ .

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The set of perturbed states is denoted  $E(\varepsilon)$ .

$$“\mathcal{J}(\varepsilon)” = L(\varepsilon, u_\varepsilon, \psi) = L(\varepsilon, u_0, \psi) + \int_0^1 \partial_u L(\varepsilon, su_\varepsilon + (1-s)u_0, \psi)(u_\varepsilon - u_0) ds$$

### Definition (Averaged adjoint state)

Given  $\varepsilon \in [0, \tau]$  and  $(u_0, u_\varepsilon) \in E(0) \times E(\varepsilon)$ , the averaged adjoint state equation is defined as follows: find  $q_\varepsilon \in X$ , such that

$$\int_0^1 \partial_u L(\varepsilon, su_\varepsilon + (1-s)u_0, q_\varepsilon)(\varphi) ds = 0 \quad \text{for all } \varphi \in X.$$

The set of averaged adjoint states is denoted  $Y(\varepsilon, u_0, u_\varepsilon)$ . We set  $Y(0, u_0) := Y(0, u_0, u_0)$ .

Thus, “ $\mathcal{J}(\varepsilon)$ ” =  $L(\varepsilon, u_\varepsilon, \psi) = L(\varepsilon, u_\varepsilon, q_\varepsilon) = L(\varepsilon, u_0, q_\varepsilon)$ .



## Averaged adjoint theorem

### Theorem (Delfour/Sturm)

Let  $\ell : \mathbf{R} \rightarrow \mathbf{R}$  be such that  $\ell(0) = 0$ . Suppose the following conditions are satisfied.

- (H1) The set of perturbed states and averaged adjoint states is non-empty for all  $\varepsilon \in [0, \tau]$ .
- (H2) For all  $u_0 \in E(0)$  and  $q_0 \in Y(0, u_0)$  the limit

$$\partial_\ell L(0, u_0, q_0) := \lim_{\varepsilon \searrow 0} \frac{L(\varepsilon, u_0, q_0) - L(0, u_0, q_0)}{\ell(\varepsilon)} \quad \text{exists.}$$

- (H3) The limit

$$R := \lim_{\varepsilon \searrow 0} \frac{L(\varepsilon, u_0, q_\varepsilon) - L(\varepsilon, u_0, q_0)}{\ell(\varepsilon)} \quad \text{exists.}$$

Then we have with  $g(\varepsilon) := L(\varepsilon, u_\varepsilon, 0)$ ,

$$d_\ell g(0) = \partial_\ell L(0, u_0, q_0) + R.$$

### Remark

For optimal control problems we usually choose  $\ell(\varepsilon) = \varepsilon$ .

## Example

Recall the example  $X = Y = H_0^1(\Omega)$ ,  $\ell(\varepsilon) := \varepsilon$ , and

$$L(\varepsilon, u, q) := \int_{\Omega} u^2 dx + \int_{\Omega} \nabla u \cdot \nabla q - \varepsilon f q dx.$$

State and averaged adjoint state equation:

$$u_{\varepsilon} \in H_0^1(\Omega), \quad \int_{\Omega} \nabla u_{\varepsilon} \cdot \nabla \varphi dx = \int_{\Omega} \varepsilon f \varphi dx \quad \text{for all } \varphi \in H_0^1(\Omega)$$

$$q_{\varepsilon} \in H_0^1(\Omega), \quad \int_{\Omega} \nabla \psi \cdot \nabla q_{\varepsilon} dx = - \int_{\Omega} (u_{\varepsilon} + u) \psi dx \quad \text{for all } \psi \in H_0^1(\Omega).$$

We have ( $u := u_0, q := q_0$ )

$$\partial_{\varepsilon} L(0, u, q) = \lim_{\varepsilon \searrow 0} \frac{L(\varepsilon, u, q) - L(0, u, q)}{\varepsilon} = - \int_{\Omega} f q dx$$

$$R = \lim_{\varepsilon \searrow 0} \frac{L(\varepsilon, u, q_{\varepsilon}) - L(\varepsilon, u, q)}{\varepsilon} = 0.$$

So  $d_{\varepsilon} g(0) = \frac{d}{d\varepsilon} L(\varepsilon, u_{\varepsilon}, 0) = - \int_{\Omega} f q dx$ .

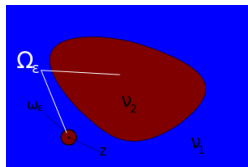
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## Application to the model problem

From now on we fix:

- an open and bounded set  $\omega \subset \mathbf{R}^d$  with  $0 \in \omega$ ,
- an open set  $\Omega \Subset D$  and a point  $z \in D \setminus \bar{\Omega}$ ,
- the perturbation  $\Omega_\varepsilon := \Omega \cup \omega_\varepsilon(z)$ , where  $\omega_\varepsilon(z) := z + \varepsilon\omega$  and  $\varepsilon \in [0, \tau]$ ,  $\tau > 0$ .



Let  $X = Y = H_0^1(D)$  and introduce the Lagrangian  $G : [0, \tau] \times X \times Y \rightarrow \mathbf{R}$  associated with the perturbation  $\Omega_\varepsilon$  by

$$L(\varepsilon, u, v) := \int_{\Gamma_0} |\mathbf{curl} u \cdot \mathbf{n} - B_d^n|^2 dS_x + \int_D v_\varepsilon(x, |\mathbf{curl} u|) \mathbf{curl} u \cdot \mathbf{curl} v dx - \langle \mathbf{F}, v \rangle.$$

where we use the abbreviation

$$v_\varepsilon(x, s) := v_1(s) \chi_{D \setminus \Omega_\varepsilon}(x) + v_2(s) \chi_{\Omega_\varepsilon}(x)$$

We will apply the averaged adjoint theorem with  $\ell(\varepsilon) = |\omega_\varepsilon|$ .

## Bound for the perturbed state

The perturbed state equation reads: find  $u_\varepsilon \in V$  such that

$$\partial_q L(\varepsilon, u_\varepsilon, 0)(v) = 0 \quad \text{for all } v \in \mathbf{V},$$

or equivalently  $u_\varepsilon \in H(D, \mathbf{curl})$  satisfies,

$$\int_D \nu_\varepsilon(x, |\mathbf{curl} u_\varepsilon|) \mathbf{curl} u_\varepsilon \cdot \mathbf{curl} v dx = \langle \mathbf{F}, v \rangle \quad \text{for all } v \in \mathbf{V}.$$

### Lemma

*There is a constant  $C > 0$ , such that for all small  $\varepsilon > 0$ ,*

$$\|u_\varepsilon - u\|_{L^2(D)^3} + \|\mathbf{curl}(u_\varepsilon - u)\|_{L^2(D)^3} \leq C\varepsilon^{d/2}.$$

## Bound for the averaged adjoint state

For  $x \in D$  and  $v \in \mathbf{R}^3$ , let  $F_\varepsilon(x, v) := \nu_\varepsilon(x, |v|)v$ . We introduce for  $\varepsilon \in [0, \tau]$  and  $v, w \in \mathbf{R}^3$ ,  $v \neq 0$ ,

$$b_\varepsilon(x, v, w) := \partial_v F_\varepsilon(x, v)w = \left( \nu_\varepsilon(x, |v|)I + \frac{\nu'_\varepsilon(x, |v|)}{|v|} v \otimes v \right) w.$$

The averaged adjoint  $q_\varepsilon \in \mathbf{V}$  is defined by

$$\int_0^1 \partial_u L(\varepsilon, su_\varepsilon + (1-s)u_0, q_\varepsilon)(v) ds = 0 \quad \text{for all } v \in \mathbf{V}.$$

This is equivalent to

$$\begin{aligned} \int_D \int_0^1 b_\varepsilon(x, \mathbf{curl}(su_\varepsilon + (1-s)u_0), \mathbf{curl}(v)) ds \cdot \mathbf{curl}(q_\varepsilon) dx \\ = - \int_{\Gamma_0} (\mathbf{curl}(u_\varepsilon + u_0) \cdot \mathbf{n} - 2B_d^n) \mathbf{curl}(v) \cdot \mathbf{n} dS_x \end{aligned}$$

for all  $v \in \mathbf{V}$ .

### Lemma

*There is a constant  $C > 0$ , such that for all small  $\varepsilon > 0$ ,*

$$\|q_\varepsilon - q\|_{H(D, \mathbf{curl})} \leq C \left( \|\mathbf{curl}(u_\varepsilon - u)\|_{L_2(D)^3} + \varepsilon^{d/2} \right).$$

# Topological Derivative

By the previous theorem, we get that the topological derivative is given by

$$G^{1 \rightarrow 2}(z) = d_\ell g(0) = \partial_\ell L(0, u, q) + R$$

if both terms exist.

## Lemma

We have

$$\begin{aligned} \partial_\ell L(0, u, q) &= \lim_{\varepsilon \searrow 0} \frac{L(\varepsilon, u, q) - L(0, u, q)}{|\omega_\varepsilon|} \\ &= (\nu_2(|\mathbf{curl} u(z)|) - \nu_1(|\mathbf{curl} u(z)|)) \mathbf{curl} u(z) \cdot \mathbf{curl} q(z) \end{aligned}$$

## Proof.

Change of variables and using that all functions are continuous at  $z$ . □

## Topological Derivative

Let us now consider the second term

$$R = \lim_{\varepsilon \searrow 0} \frac{L(\varepsilon, u, q_\varepsilon) - L(\varepsilon, u, q)}{\ell(\varepsilon)}$$

where  $\ell(\varepsilon) = |\omega_\varepsilon|$ . First note that testing the state equation for  $\varepsilon = 0$  with  $\mathbf{v} := \mathbf{q}_\varepsilon - \mathbf{q}$  yields

$$\int_{\mathbf{D}} \nu_0(x, |\mathbf{curl} u|) \mathbf{curl} u \cdot \mathbf{curl}(\mathbf{q}_\varepsilon - \mathbf{q}) \, dx = \langle \mathbf{F}, \mathbf{q}_\varepsilon - \mathbf{q} \rangle. \quad (1)$$

Therefore

$$\begin{aligned} L(\varepsilon, u, \mathbf{q}_\varepsilon) - L(\varepsilon, u, \mathbf{q}) &= \int_{\mathbf{D}} \nu_\varepsilon(x, |\mathbf{curl} u|) \mathbf{curl} u \cdot \mathbf{curl}(\mathbf{q}_\varepsilon - \mathbf{q}) \, dx - \langle \mathbf{F}, \mathbf{q}_\varepsilon - \mathbf{q} \rangle \\ &\stackrel{(1)}{=} \int_{\mathbf{D}} (\nu_\varepsilon(x, |\mathbf{curl} u|) - \nu_0(x, |\mathbf{curl} u|)) \mathbf{curl} u \cdot \mathbf{curl}(\mathbf{q}_\varepsilon - \mathbf{q}) \, dx \\ &= \int_{\omega_\varepsilon} (\nu_2(|\mathbf{curl} u|) - \nu_1(|\mathbf{curl} u|)) \mathbf{curl} u \cdot \mathbf{curl}(\mathbf{q}_\varepsilon - \mathbf{q}) \, dx. \end{aligned}$$



## Definition

We use the following Helmholtz decomposition for  $u_\varepsilon$ :

$$u_\varepsilon = \nabla\phi_\varepsilon + w_\varepsilon, \quad \phi_\varepsilon \in H_0^1(D), \quad w_\varepsilon \in H_0^1(D)^3.$$

Similarly we decompose  $q_\varepsilon$ :

$$q_\varepsilon = \nabla\psi_\varepsilon + z_\varepsilon, \quad \psi_\varepsilon \in H_0^1(D), \quad z_\varepsilon \in H_0^1(D)^3.$$

## Definition

The variation of the averaged adjoint state  $q_\varepsilon$  is defined pointwise a.e. in  $\mathbf{R}^d$  by

$$Q_\varepsilon(x) := \frac{\tilde{z}_\varepsilon(T_\varepsilon(x)) - \tilde{z}(T_\varepsilon(x))}{\varepsilon},$$

and the variation of the direct state  $u_\varepsilon$  is defined pointwise a.e. in  $\mathbf{R}^d$  by

$$W_\varepsilon(x) := \frac{\tilde{w}_\varepsilon(T_\varepsilon(x)) - \tilde{w}(T_\varepsilon(x))}{\varepsilon}.$$

Here,  $\tilde{q}_\varepsilon$ ,  $\tilde{q}$ ,  $\tilde{u}_\varepsilon$ ,  $\tilde{u}$  are extensions of  $q_\varepsilon$ ,  $q$ ,  $u_\varepsilon$ ,  $u$  by zero.

Invoking the change of variables  $y = T_\varepsilon(x) = z + \varepsilon x$ , we get

$$\begin{aligned} & \frac{L(\varepsilon, u, q_\varepsilon) - L(\varepsilon, u, q)}{|\omega_\varepsilon|} \\ &= \frac{1}{|\omega|} \int_\omega (\nu_2(|\mathbf{curl} u(T_\varepsilon(x))|) - \nu_1(|\mathbf{curl} u(T_\varepsilon(x))|)) \mathbf{curl} u(T_\varepsilon(x)) \cdot \mathbf{curl} Q_\varepsilon \, dx. \end{aligned}$$

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Invoking the change of variables  $y = T_\varepsilon(x) = z + \varepsilon x$ , we get

$$\begin{aligned} & \frac{L(\varepsilon, u, q_\varepsilon) - L(\varepsilon, u, q)}{|\omega_\varepsilon|} \\ &= \frac{1}{|\omega|} \int_\omega (\nu_2(|\mathbf{curl} u(T_\varepsilon(x))|) - \nu_1(|\mathbf{curl} u(T_\varepsilon(x))|)) \mathbf{curl} u(T_\varepsilon(x)) \cdot \mathbf{curl} Q_\varepsilon \, dx. \\ & \xrightarrow{\varepsilon \rightarrow 0} ??? \end{aligned}$$

## Conjecture

There exists  $W \in BL(\mathbf{R}^3)$ , such that

$$\begin{aligned} \partial W_\varepsilon &\rightarrow \partial W && \text{in } L_2(\mathbf{R}^3)^{3 \times 3} \quad \text{as } \varepsilon \searrow 0 \\ \mathbf{curl}(W_\varepsilon) &\rightarrow \mathbf{curl}(W) && \text{in } L_2(\mathbf{R}^3)^3 \quad \text{as } \varepsilon \searrow 0. \end{aligned}$$

Moreover,  $W$  satisfies

$$\begin{aligned} \int_{\mathbf{R}^3} (F_\omega(x, \mathbf{curl} W + \mathbf{curl}(w_0)(z)) - F_\omega(x, \mathbf{curl} w_0(z))) \cdot \mathbf{curl} v \, dx = \\ - \int_\omega (\nu_1(|\mathbf{curl} w_0(z)|) - \nu_2(|\mathbf{curl} w_0(z)|)) \mathbf{curl} w_0(z) \cdot \mathbf{curl} v \, dx \end{aligned}$$

for all  $v \in BL(\mathbf{R}^3)$ . Here  $F_\omega(x, y) := \nu_1(y)\chi_{\mathbf{R}^3 \setminus \omega}(x) + \nu_2(y)\chi_\omega(x)$ .

## Corollary of Conjecture

The weak limit  $\mathbf{Q}$  of  $\mathbf{Q}_\varepsilon = (\mathbf{q}_\varepsilon - \mathbf{q}_0)/\varepsilon \circ T_\varepsilon$  satisfies

$$\begin{aligned} & \int_{\mathbf{R}^3} \left( \int_0^1 \partial_y \tilde{F}_\omega(x, s \mathbf{curl} W + \mathbf{curl} u_0(z)) \mathbf{curl} \bar{\mathbf{v}} \, ds \right) \cdot \mathbf{curl} \mathbf{Q} \, dx \\ &= - \int_{\mathbf{R}^3} \left[ \int_0^1 \partial_y \tilde{F}_\omega(x, s \mathbf{curl} W + \mathbf{curl} u_0(z)) \mathbf{curl}(\bar{\mathbf{v}}) \, ds \right. \\ & \quad \left. - \partial_y \tilde{F}_\omega(x, \mathbf{curl} u_0(z)) \mathbf{curl}(\bar{\mathbf{v}}) \right] \cdot \mathbf{curl} \mathbf{q}_0(z) \, dx \\ & - \int_\omega [\partial_y F_2(\mathbf{curl} u_0(z)) - \partial_y F_1(\mathbf{curl} u_0(z))] \mathbf{curl}(\bar{\mathbf{v}}) \cdot \mathbf{curl} \mathbf{q}_0(z) \, dx. \end{aligned}$$

## Topological Derivative

Suppose that Conjecture holds. Then

$$\begin{aligned} & \frac{L(\varepsilon, u, q_\varepsilon) - L(\varepsilon, u, q)}{|\omega_\varepsilon|} \\ &= \frac{1}{|\omega|} \int_\omega (\nu_2(|\mathbf{curl} u(T_\varepsilon(x))|) - \nu_1(|\mathbf{curl} u(T_\varepsilon(x))|)) \mathbf{curl} u(T_\varepsilon(x)) \cdot \mathbf{curl} Q_\varepsilon \, dx. \\ &\xrightarrow{\varepsilon \rightarrow 0} (\nu_2(|\mathbf{curl} u(z)|) - \nu_1(|\mathbf{curl} u(z)|)) \mathbf{curl} u(z) \cdot \frac{1}{|\omega|} \int_\omega \mathbf{curl} Q \, dx \end{aligned}$$

For the topological derivative, we would get

$$\partial J(z) = d_\ell g(0) = \partial_\ell L(0, u_0, q_0) + R$$

$$= (\nu_2(|\mathbf{curl} u(z)|) - \nu_1(|\mathbf{curl} u(z)|)) \mathbf{curl} u(z) \cdot \left( \mathbf{curl} q(z) + \frac{1}{|\omega|} \int_\omega \mathbf{curl} Q \, dx \right)$$

## It's not the end of the story

- computation of  $Q$  is numerically infeasible
- BUT: we can eliminate it by testing the equation for  $W$  by  $Q$  and using the fundamental theorem of calculus.

$$\begin{aligned} R(u, p) &= (\nu_1(|\mathbf{curl} u(z)|) - \nu_2(|\mathbf{curl} u(z)|)) \mathbf{curl} u(z) \cdot \frac{1}{|\omega|} \int_{\omega} \mathbf{curl} Q \, dx \\ &= \frac{1}{|\omega|} \int_{\mathbf{R}^3} \left[ \tilde{F}_{\omega}(x, \mathbf{curl} W + \mathbf{curl} w_0(z)) - \tilde{F}_{\omega}(x, \mathbf{curl} w_0(z)) \right. \\ &\quad \left. - \partial_y \tilde{F}_{\omega}(x, \mathbf{curl} w_0(z)) \mathbf{curl}(W) \right] \cdot \mathbf{curl} q_0 \, dx \\ &\quad + \frac{1}{|\omega|} \int_{\omega} [\partial_y F_2(\mathbf{curl} w_0(z)) - \partial_y F_1(\mathbf{curl} w_0(z))] \mathbf{curl}(W) \cdot \mathbf{curl} q_0(z) \, dx. \end{aligned}$$