

# A second order multipoint flux mixed finite element method on hybrid meshes

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## Porous media modeling

Model equations for single-phase flow:

Conservation of mass

$$\operatorname{div} \mathbf{u} = f$$

Darcy's law

$$\mathbf{u} = -K \nabla p$$

Quantity of interest :  $p$

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Second order form

$$\begin{aligned} -\operatorname{div} (K \nabla p) &= f && \text{in } \Omega \\ p &= 0 && \text{on } \partial\Omega. \end{aligned}$$

- (i) Discontinuous schemes (DFVM), (DG) for local mass conservation
- (ii) Not accurate for rough coefficients (local arithmetic averaging of  $K$ )

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- (i) Discontinuous schemes (DFVM), (DG) for local mass conservation
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Mixed form

$$\begin{aligned} K^{-1} \mathbf{u} + \nabla p &= 0 && \text{in } \Omega \\ \operatorname{div} \mathbf{u} &= f && \text{in } \Omega \\ p &= 0 && \text{on } \partial\Omega. \end{aligned}$$

- (i) Handles rough coefficients better (local **harmonic averaging** of  $K$ )
- (ii) Have to solve a full saddle point problem... or do you ?  $\Rightarrow$  **MFMFE**

## Variational formulation

$$\begin{aligned}K^{-1}\mathbf{u} + \nabla p &= 0 && \text{in } \Omega \\ \operatorname{div} \mathbf{u} &= f && \text{in } \Omega \\ p &= 0 && \text{on } \partial\Omega.\end{aligned}$$

### Variational formulation

$$\begin{aligned}(K^{-1}\mathbf{u}, \mathbf{v}) - (p, \operatorname{div} \mathbf{v}) &= 0 && \forall \mathbf{v} \in H(\operatorname{div}, \Omega) \\ (\operatorname{div} \mathbf{u}, q) &= (f, q) && \forall q \in L^2(\Omega)\end{aligned}$$

## Discrete variational formulation

$$\begin{aligned}K^{-1}\mathbf{u} + \nabla p &= 0 && \text{in } \Omega \\ \operatorname{div} \mathbf{u} &= f && \text{in } \Omega \\ p &= 0 && \text{on } \partial\Omega.\end{aligned}$$

Discrete variational formulation

$$\begin{aligned}(K^{-1}\mathbf{u}_h, \mathbf{v}_h) - (p_h, \operatorname{div} \mathbf{v}_h) &= 0 && \forall \mathbf{v}_h \in \mathbf{V}_h \subseteq H(\operatorname{div}, \Omega) \\ (\operatorname{div} \mathbf{u}_h, q_h) &= (f, q_h) && \forall q_h \in Q_h \subseteq L^2(\Omega)\end{aligned}$$

Problem : we have to solve a full (indefinite) saddle point system ...

## Mass lumping

$$\begin{aligned}K^{-1}\mathbf{u} + \nabla p &= 0 && \text{in } \Omega \\ \operatorname{div} \mathbf{u} &= f && \text{in } \Omega \\ p &= 0 && \text{on } \partial\Omega.\end{aligned}$$

Discrete variational formulation via *mass lumping* (MFMFE)

$$\begin{aligned}(K^{-1}\mathbf{u}_h, \mathbf{v}_h)_h - (p_h, \operatorname{div} \mathbf{v}_h) &= 0 && \forall \mathbf{v}_h \in \mathbf{V}_h \subseteq H(\operatorname{div}, \Omega) \\ (\operatorname{div} \mathbf{u}_h, q_h) &= (f, q_h) && \forall q_h \in Q_h \subseteq L^2(\Omega)\end{aligned}$$

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For appropriate spaces  $\mathbf{V}_h$ ,  $Q_h$  and  $(\cdot, \cdot)_h$ , the *lumped mass matrix*  $M_h$  is block-diagonal, and the variable  $\mathbf{u}_h$  can be eliminated efficiently.

$$\begin{pmatrix} M_h & -C^\top \\ C & 0 \end{pmatrix} \begin{pmatrix} \mathbf{u}_h \\ p_h \end{pmatrix} = \begin{pmatrix} 0 \\ \mathbf{f} \end{pmatrix} \quad \Longrightarrow \quad CM_h^{-1}C^\top p_h = \mathbf{f}$$

The problem reduces to symmetric, positive definite cell-centered system for the pressure (CCFD)



# Discretization

Discrete variational formulation via *mass lumping* (MFMFE)

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M. Wheeler, I. Yotov *A multipoint flux mixed finite element method*. SIAM 2006

$$V(T) = \text{BDM}_1(T) := P_1(T)^2$$

$$Q(T) = P_0(T)$$

$$(\mathbf{u}_h, \mathbf{v}_h)_h := \frac{|T|}{3} \sum_{i=1}^3 \mathbf{u}_h(r_i) \mathbf{v}_h(r_i)$$

$r_i$  vertex

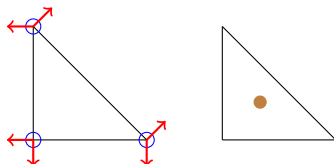


Figure: DOFs of  $V(T)$  (left) and  $Q(T)$  (right). Blue circles are quadrature points.

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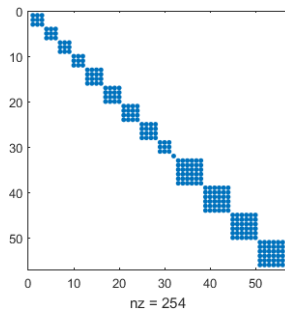
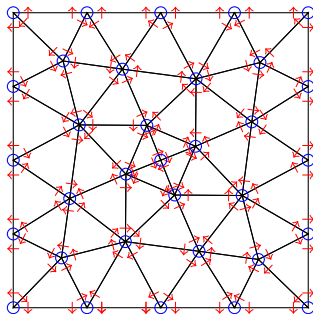
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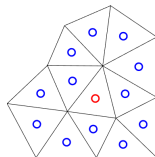
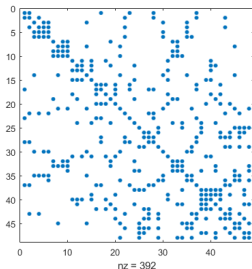


Figure: Matrix  $CM_h^{-1}C^T$  (left), stencil of the method (right)

## Convergence analysis

Summary of the convergence results

$$\|\mathbf{u} - \mathbf{u}_h\| = O(h) \quad \text{and} \quad \|\pi_h^0 p - p_h\| = O(h^2)$$

Relevant properties

- (i)  $P_0(T)^2 \subseteq \mathbf{V}(T)$  and  $P_0(T) \subseteq Q(T)$  such that  $\operatorname{div} \mathbf{V}(T) \subseteq Q(T)$
- (ii) The quadrature rule is exact for  $P_0(T)^2 \times \mathbf{V}(T)$
- (iii) The quadrature rule induces a norm on  $\mathbf{V}(T)$

## Convergence analysis

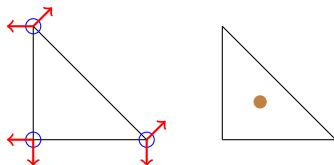
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- (ii) The quadrature rule is exact for  $P_0(T)^2 \times \mathbf{V}(T)$  ✓
- (iii) The quadrature rule induces a norm on  $\mathbf{V}(T)$  ✓

**Wheeler-Yotov element** :  $\mathbf{V}(T) = \text{BDM}_1(T) = P_1(T)^2$



**Figure:** DOFs of  $\mathbf{V}(T)$  (left) and  $Q(T)$  (right). Blue circles are quadrature points. The quadrature rule is exact for  $P_1(T)$ .

## Higher order candidates

Natural extension of the first order estimates

$$\|\mathbf{u} - \mathbf{u}_h\| = O(h^2) \quad \text{and} \quad \|\pi_h^1 p - p_h\| = O(h^3)$$

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- (ii) The quadrature rule is exact for  $P_1(T)^2 \times \mathbf{V}(T)$  ✗
- (iii) The quadrature rule induces a norm on  $\mathbf{V}(T)$  ✗

**First candidate** :  $\mathbf{V}(T) = \text{BDM}_2(T) = P_2(T)^2$

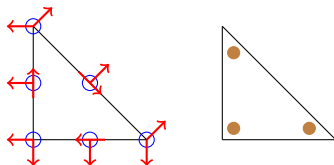


Figure: DOFs of  $\mathbf{V}(T)$  (left) and  $Q(T)$  (right). Blue circles are quadrature points.

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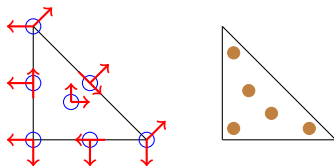
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- (i)  $P_1(T)^2 \subseteq \mathbf{V}(T)$  and  $P_1(T) \subseteq Q(T)$  such that  $\text{div } \mathbf{V}(T) \subseteq Q(T)$  ✓
- (ii) The quadrature rule is exact for  $P_1(T)^2 \times \mathbf{V}(T)$  ✓
- (iii) The quadrature rule induces a norm on  $\mathbf{V}(T)$  ✓

**Second candidate** :  $\mathbf{V}(T) = \text{BDM}_2^+(T) = P_2(T)^2 \oplus b_3 \cdot [1, 0]^T \oplus b_3 \cdot [0, 1]^T$



**Figure:** DOFs of  $\mathbf{V}(T)$  (left) and  $Q(T)$  (right). Blue circles are quadrature points. The quadrature rule is exact for  $P_3(T) \oplus b_3 \cdot P_1(T)$ .



## Higher order candidates

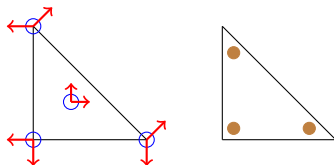
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- (iii) The quadrature rule induces a norm on  $\mathbf{V}(T)$  ✓

**Third candidate** :  $\mathbf{V}(T) = \text{RT}_1(T) := P_1(T)^2 + \mathbf{x} \cdot P_1^h(T)$



**Figure:** DOFs of  $\mathbf{V}(T)$  (left) and  $Q(T)$  (right). Blue circles are quadrature points. The quadrature rule is exact for  $P_2(T)$ .

## A new theory

Split the error in  $\|\mathbf{u} - \mathbf{u}_h\|_{L^2(\Omega)} \leq \|\mathbf{u} - \Pi_h \mathbf{u}\|_{L^2(\Omega)} + \|\Pi_h \mathbf{u} - \mathbf{u}_h\|_{L^2(\Omega)}$

$$\begin{aligned}(\Pi_h \mathbf{u} - \mathbf{u}_h, \mathbf{v}_h)_h - (\pi_h^1 p - p_h, \operatorname{div} \mathbf{v}_h) &= (\Pi_h \mathbf{u} - \mathbf{u}, \mathbf{v}_h) + \sigma_h(\Pi_h \mathbf{u}, \mathbf{v}_h) \\ (\operatorname{div}(\Pi_h \mathbf{u} - \mathbf{u}_h), q_h) &= 0\end{aligned}$$

with  $\sigma_h(\Pi_h \mathbf{u}, \mathbf{v}_h) = (\Pi_h \mathbf{u}, \mathbf{v}_h)_h - (\Pi_h \mathbf{u}, \mathbf{v}_h)$ .

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- (I)  $\operatorname{div}(\Pi_h \mathbf{u} - \mathbf{u}_h) = 0 \Rightarrow \Pi_h \mathbf{u} - \mathbf{u}_h \in P_1(T)^2$
- (II)  $\sigma_h(\mathbf{u}_h, \mathbf{v}_h) = 0$  if  $\mathbf{u}_h, \mathbf{v}_h \in P_1(T)^2$

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with  $\sigma_h(\Pi_h \mathbf{u}, \mathbf{v}_h) = (\Pi_h \mathbf{u}, \mathbf{v}_h)_h - (\Pi_h \mathbf{u}, \mathbf{v}_h)$ .

$$(I) \operatorname{div}(\Pi_h \mathbf{u} - \mathbf{u}_h) = 0 \quad \Rightarrow \quad \Pi_h \mathbf{u} - \mathbf{u}_h \in P_1(T)^2$$

$$(II) \sigma_h(\mathbf{u}_h, \mathbf{v}_h) = 0 \text{ if } \mathbf{u}_h, \mathbf{v}_h \in P_1(T)^2$$

Taking  $\mathbf{v}_h = \Pi_h \mathbf{u} - \mathbf{u}_h$  and  $q_h = \pi_h^1 p - p_h$ , we obtain

$$\begin{aligned}\|\Pi_h \mathbf{u} - \mathbf{u}_h\|_h^2 &= (\Pi_h \mathbf{u} - \mathbf{u}, \Pi_h \mathbf{u} - \mathbf{u}_h) + \sigma_h(\Pi_h \mathbf{u}, \Pi_h \mathbf{u} - \mathbf{u}_h) \\&= (\Pi_h \mathbf{u} - \mathbf{u}, \Pi_h \mathbf{u} - \mathbf{u}_h) + \sigma_h(\Pi_h \mathbf{u} - \pi_h^1 \mathbf{u}, \Pi_h \mathbf{u} - \mathbf{u}_h) \\&\leq \|\Pi_h \mathbf{u} - \mathbf{u}\|_0 \|\Pi_h \mathbf{u} - \mathbf{u}_h\|_0 + c \|\Pi_h \mathbf{u} - \pi_h^1 \mathbf{u}\|_0 \|\Pi_h \mathbf{u} - \mathbf{u}_h\|_0 \\&\leq Ch^2 \|\mathbf{u}\|_{H^2(\mathcal{T}_h)} \|\Pi_h \mathbf{u} - \mathbf{u}_h\|_0\end{aligned}$$

## A new theory

### Theorem

$$\|\mathbf{u} - \mathbf{u}_h\| = O(h^2) \quad \text{and} \quad \|\pi_h^0(p - p_h)\| = O(h^3)$$

#### Relevant properties

- (i)  $P_1(T)^2 \subset \mathbf{V}(T)$  and  $P_1(T) \subset Q(T)$  such that  $\operatorname{div} \mathbf{V}(T) \subseteq Q(T)$  ✓
- (ii<sub>a</sub>)  $\exists \tilde{\mathbf{V}}(T) \subset \mathbf{V}(T)$  s.t.  $\mathbf{v} \in \mathbf{V}(T)$  with  $\operatorname{div} \mathbf{v} \in \operatorname{div} \tilde{\mathbf{V}}(T)$  imply  $\mathbf{v} \in \tilde{\mathbf{V}}(T)$  ✓
- (ii<sub>b</sub>) The quadrature rule is exact for  $P_1(T)^2 \times \tilde{\mathbf{V}}(T)$  ✓
- (iii) The quadrature rule induces a norm on  $\mathbf{V}(T)$  ✓

**Third candidate** :  $\mathbf{V}(T) = \mathbf{RT}_1(T) := P_1(T)^2 + \mathbf{x} \cdot P_1^h(T)$

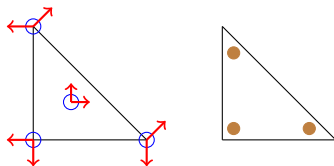


Figure: DOFs of  $V(T)$  (left) and  $Q(T)$  (right). Blue circles are quadrature points.

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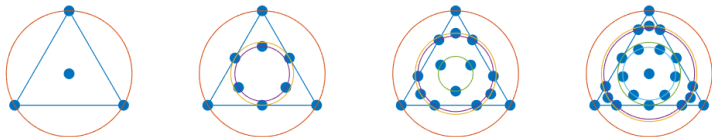
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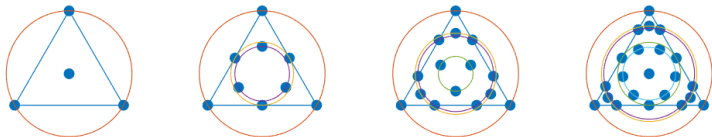


Similar concept in the paper by Geervers, et al, 2018



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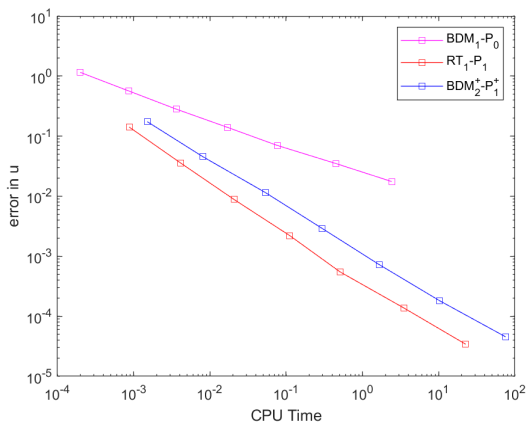


Similar concept in the paper by Geevers, et al, 2018

- (iv) Application to wave propagation









$$\begin{aligned} \partial_t \mathbf{u} + \nabla p &= f && \text{in } \Omega \\ \partial_t p + \operatorname{div} \mathbf{u} &= g && \text{in } \Omega \\ p &= 0 && \text{on } \partial\Omega. \end{aligned}$$

## Comparison



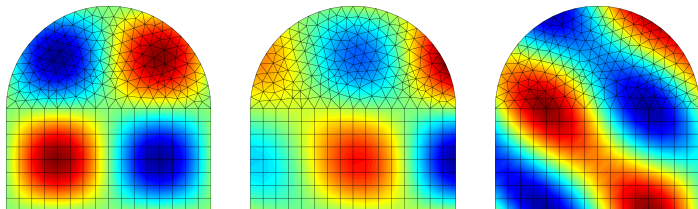
The  $RT_1 - P_1$  pair is about 3x faster than the  $BDM_2^+ - P_1^+$  pair.

# Hybrid meshes

	$\dim \mathbf{V}(T)$	$\dim Q(T)$	$\ \mathbf{u} - \mathbf{u}_h\ _0$	$\ \pi_h^0 p - p_h\ _0$	DOFs for $\mathbf{u}_h$	DOFs for $p_h$
$\text{BDM}_1 - \text{P}_0$	6+0	1	$O(h)$	$O(h^2)$		
$\text{RT}_1 - \text{P}_1$	6+2	3	$O(h^2)$	$O(h^3)$		
$\text{BDM}_1 - \text{P}_0$	8+0	1	$O(h)$	$O(h^2)$		
$\text{BDFM}_2 - \text{P}_1$	8+2	3	$O(h^2)$	$O(h^3)$		

## Numerical tests

$$p = \sin(\pi x) \sin(\pi y) \quad K = \begin{pmatrix} 4 + (x + 2)^2 + y^2 & 1 + \sin(xy) \\ 1 + \sin(xy) & 2 \end{pmatrix}$$



**Figure:** Snapshots of the pressure  $p_h$  (left) and the two velocity components  $u_{x,h}$ ,  $u_{y,h}$  (middle, right) for the second order approximation.

$h$	DOF $u$	DOF $p$	$\ u - u_h\ $	eoc	$\ \pi_h^0(p - p_h)\ $	eoc
$2^{-1}$	164	84	0.078309	—	0.033106	—
$2^{-2}$	724	396	0.013097	2.57	0.002864	3.53
$2^{-3}$	2498	1386	0.002275	2.52	0.000391	2.87
$2^{-4}$	9738	5466	0.000484	2.23	0.000049	2.99
$2^{-5}$	40230	22770	0.000099	2.28	0.000005	3.13

**Table:** Degrees of freedom, relative discretization errors, and convergence rates for the second order multipoint flux finite element method.

## Summary

- Introduced the multipoint flux mixed finite element method (MFMFE)
- Presented the first order approximation introduced by Wheeler and Yotov
- Proposed an extension to second order approximations



H. Egger, B. Radu *A second order multipoint flux mixed finite element method on hybrid meshes*, TU Darmstadt, 12/2018 arXiv: 1812.03938

### A few additional remarks

- Extension to the  $3D$  case has also been done.
- The framework can be used to design even higher order approximations
- We can devise local post-processing strategies for the pressure
- The techniques can also be applied for the wave and Maxwell's equations

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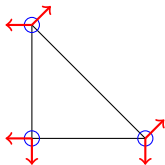
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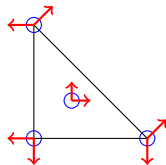
## Thank you for your attention

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## Triangles

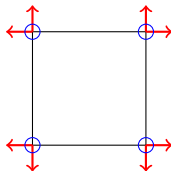


(a)  $BDM_1$  first order element

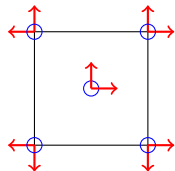


(b)  $RT_1$  second order element

## Parallelograms

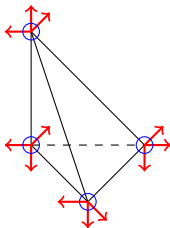


(a)  $BDM_1$  first order element

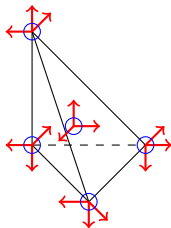


(b)  $BDFM_2$  second order element

## Tetrahedra

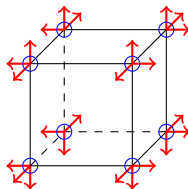


(a)  $BDM_1$  first order element

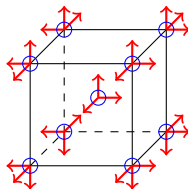


(b)  $RT_1$  second order element

## Hexahedra



(a)  $eBDDF_1$  first order element



(b) ? - second order element