

# Simulating a Heart Valve using a Varying Permeability Approach

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## GOAL

Patient-specific model of total heart function

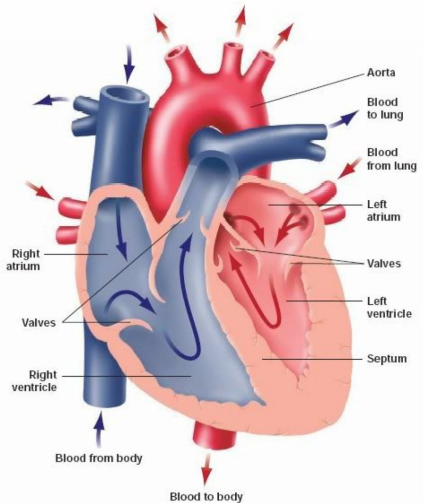
Physics involved:

- Electrophysiology
- Mechanics
- Fluid Dynamics

Code basis:



Cardiac Arrhythmia Research Package



<http://getdrawings.com/>

## GOAL

Model the impact of a valve in the human heart

## ISSUES

- complex multiphysics problem
- interface conditions
- projecting boundary data from one domain into the other
- costly in terms of computation time

## IDEA

Penalize the Navier-Stokes equations to model obstacles in the fluid domain (fictitious domain approach<sup>1</sup>)

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<sup>1</sup>Khodor Khadra et al. “Fictitious domain approach for numerical modelling of Navier–Stokes equations”. In: *International Journal for Numerical Methods in Fluids* 34.8 (2000), pp. 651–684.

## BENEFITS

- flexible and robust computational technique
- straight forward implementation
- saving computation time
- different media are accounted for by their physical characteristics assigned to the grid points
- time-dependent geometry and material properties can easily be taken into account



# Navier-Stokes-Brinkman (NSB) Model

Domain and Boundary:

$$\begin{aligned}\Omega &= \Omega_f \cup \Omega_p \cup \Omega_s \\ \Gamma &= \Gamma_e \cup \Gamma_n\end{aligned}\quad (1)$$

Navier-Stokes-Brinkmann Model:

$$\begin{cases} \rho (\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} - \nabla \cdot \boldsymbol{\sigma}(\mathbf{u}, p)) + \frac{\mu}{K} \mathbf{u} = \rho \mathbf{f} & \text{in } \mathbb{R}^+ \times \Omega \\ \nabla \cdot \mathbf{u} = 0, & \text{in } \mathbb{R}^+ \times \Omega \\ \mathbf{u}(0, \cdot) = \mathbf{u}_0 & \text{in } \Omega \\ \mathbf{u} = \mathbf{g}, & \text{on } \Gamma_e \\ \mathbf{n} \cdot \boldsymbol{\sigma}(\mathbf{u}, p) = \mathbf{h}, & \text{on } \Gamma_n \end{cases} \quad (2)$$

Stress Tensor  $\boldsymbol{\sigma}$ :

$$\begin{aligned}\boldsymbol{\sigma}(\mathbf{u}, p) &= -p \mathbb{1} + 2\mu \boldsymbol{\varepsilon}(\mathbf{u}) \\ \boldsymbol{\varepsilon}(\mathbf{u}) &= \frac{1}{2}(\nabla \mathbf{u} + \nabla \mathbf{u}^\top)\end{aligned}\quad (3)$$

$$K(t, \mathbf{x}) = \begin{cases} K_f \rightarrow +\infty, & \forall \mathbf{x} \in \Omega_f \\ K_p, & \forall \mathbf{x} \in \Omega_p \\ K_s \rightarrow 0^+, & \forall \mathbf{x} \in \Omega_s \end{cases} \quad (4)$$

- $\Omega_f$ :  $\frac{\mu}{K} \mathbf{u} \rightarrow 0$  classical Navier–Stokes equations
- $\Omega_p$ : full Navier-Stokes-Brinkmann equations
- $\Omega_s$ :  $\frac{\mu}{K} \mathbf{u}$  forces  $\mathbf{u} \rightarrow 0$  Darcy's law (5)

$$\nabla p + \frac{\mu}{K} \mathbf{u} = \rho \mathbf{f} \quad (5)$$

**Note:** No-slip condition on  $\partial\Omega_s$  is asymptotically satisfied

# Variational Arbitrary Lagrangian Eulerian (ALE) Formulation

The Navier-Stokes-Brinkman Model for a moving domain :

Find  $\mathbf{u} \in \mathcal{S}_u$  and  $p \in \mathcal{S}_p$ , such that  $\forall \mathbf{w} \in \mathcal{V}_u$  and  $q \in \mathcal{V}_p$ :

$$\begin{aligned} & \int_{\Omega} \mathbf{w} \cdot \rho \left( \partial_t \mathbf{u} + (\mathbf{u} - \hat{\mathbf{u}}) \cdot \nabla \mathbf{u} - \mathbf{f} + \frac{\nu}{K} \mathbf{u} \right) d\Omega \\ & + \int_{\Omega} \boldsymbol{\varepsilon}(\mathbf{w}) : \boldsymbol{\sigma}(\mathbf{u}, p) d\Omega - \int_{\Gamma_n} \mathbf{w} \cdot \mathbf{h} d\Gamma + \int_{\Omega} q \nabla \cdot \mathbf{u} d\Omega = 0 \end{aligned} \quad (6)$$

$\mathbf{u}$  ... fluid velocity

$\hat{\mathbf{u}}$  ... fluid domain velocity respective to a reference domain

# Residual-Based Variational Multiscale Formulation

RBVMS<sup>2</sup> formulation of the ALE NSB equations:

Decomposition into coarse ( $^h$ ) and fine ( $'$ ) scale subspaces:

$$\begin{aligned} \mathcal{S}_u &= \mathcal{S}_u^h \oplus \mathcal{S}'_u & \mathcal{S}_p &= \mathcal{S}_p^h \oplus \mathcal{S}'_p \\ \mathcal{V}_u &= \mathcal{V}_u^h \oplus \mathcal{V}'_u & \mathcal{V}_p &= \mathcal{V}_p^h \oplus \mathcal{V}'_p \end{aligned} \quad (7)$$

Model the fine scale velocity and the fine scale pressure:

$$\begin{aligned} \mathbf{u}' &= -\frac{\tau_{SUPS}}{\rho} \mathbf{r}^M(\mathbf{u}^h, p^h) \\ p' &= -\rho \nu_{LSIC} r^C(\mathbf{u}^h) \end{aligned} \quad (8)$$

Residuals  $\mathbf{r}^M$  and  $r^C$ :

$$\begin{aligned} \mathbf{r}^M(\mathbf{u}^h, p^h) &= \rho \left( \partial_t \mathbf{u}^h + \left( \mathbf{u} - \hat{\mathbf{u}} \right) \cdot \nabla \mathbf{u} - \mathbf{f}^h + \frac{\nu}{K} \mathbf{u}^h \right) - \nabla \cdot \boldsymbol{\sigma}(\mathbf{u}^h, p^h) \\ r^C(\mathbf{u}^h) &= \nabla \cdot \mathbf{u}^h \end{aligned} \quad (9)$$

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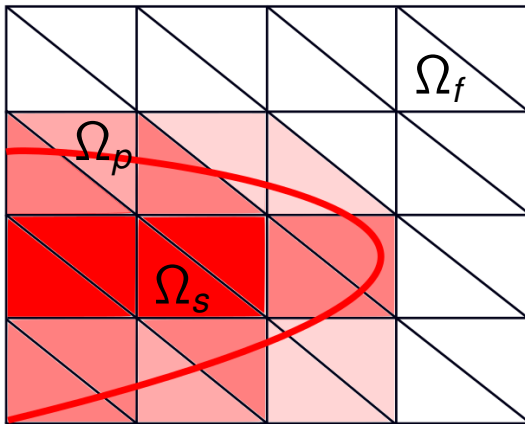
<sup>2</sup>Yuri Bazilevs, Kenji Takizawa, and Tayfun E. Tezduyar. *Computational Fluid-Structure Interaction: Methods and Applications*. English. John Wiley and Sons, 2013.

# Residual-Based Variational Multiscale Formulation

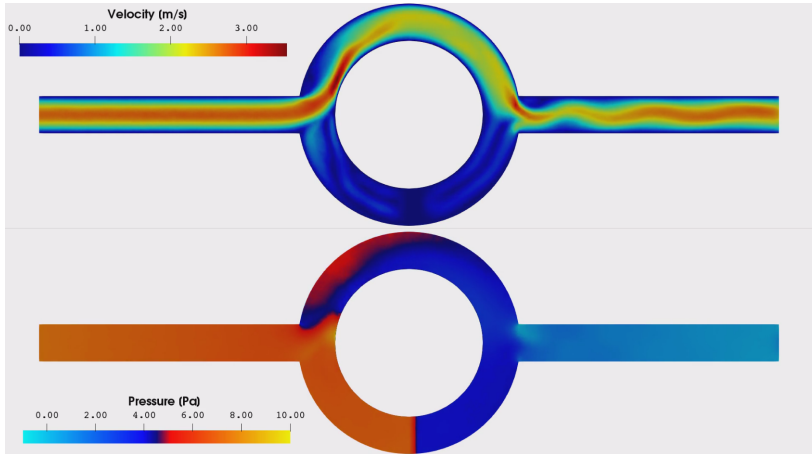
Find  $\mathbf{u}^h \in \mathcal{S}_u^h$  and  $p^h \in \mathcal{S}_p^h$ , such that  $\forall \mathbf{w}^h \in \mathcal{V}_u^h$  and  $q^h \in \mathcal{V}_p^h$ :

$$\begin{aligned}
 & \int_{\Omega} \mathbf{w}^h \cdot \rho \left( \partial_t \mathbf{u}^h + \left( \mathbf{u}^h - \hat{\mathbf{u}}^h \right) \cdot \nabla \mathbf{u}^h - \mathbf{f}^h + \frac{\nu}{K} \mathbf{u}^h \right) d\Omega + \int_{\Omega} \boldsymbol{\varepsilon}(\mathbf{w}^h) : \boldsymbol{\sigma}(\mathbf{u}^h, p^h) d\Omega \\
 & - \int_{\Gamma_n} \mathbf{w}^h \cdot \mathbf{h}^h d\Gamma + \int_{\Omega_t} q^h \nabla \cdot \mathbf{u}^h d\Omega \\
 & + \sum_{e=1}^{n_{el}} \int_{\Omega_e} \tau_{SUPS} \left( \left( \mathbf{u}^h - \hat{\mathbf{u}}^h \right) \cdot \nabla \mathbf{w}^h + \frac{\nabla q^h}{\rho} - \frac{\nu}{K} \mathbf{w}^h \right) \cdot \mathbf{r}^M(\mathbf{u}^h, p^h) d\Omega \\
 & + \sum_{e=1}^{n_{el}} \int_{\Omega_e} \rho \nu_{LSIC} \nabla \cdot \mathbf{w}^h \mathbf{r}^C(\mathbf{u}^h) d\Omega \\
 & - \sum_{e=1}^{n_{el}} \int_{\Omega_e} \tau_{SUPS} \mathbf{w}^h \cdot \left( \mathbf{r}^M(\mathbf{u}^h, p^h) \cdot \nabla \mathbf{u}^h \right) d\Omega \\
 & - \sum_{e=1}^{n_{el}} \int_{\Omega_e} \frac{\nabla \mathbf{w}^h}{\rho} : \left( \tau_{SUPS} \mathbf{r}^M(\mathbf{u}^h, p^h) \right) \otimes \left( \tau_{SUPS} (\mathbf{r}^M(\mathbf{u}^h, p^h)) \right) d\Omega = 0
 \end{aligned} \tag{10}$$

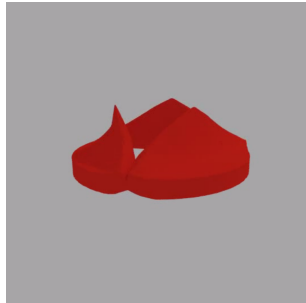
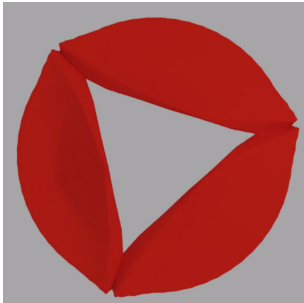
How to choose the permeability parameters to model the valve?



# Proof of Concept



## 1st STEP: Moving Valve



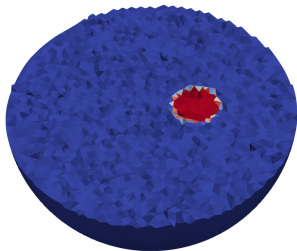
Adapt the penalization term to fulfill no-slip condition<sup>3</sup> :

$$\frac{\mu}{K} \mathbf{u} \rightarrow \frac{\mu}{K} (\mathbf{u} - \mathbf{u}_S) \quad (11)$$

<sup>3</sup>Dmitry Kolomenskiy and Kai Schneider. “A Fourier spectral method for the Navier-Stokes equations with volume penalization for moving solid obstacles”. In: *J. Comput. Physics* 228 (2009), pp. 5687–5709.

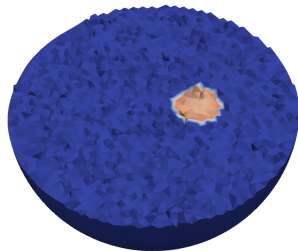


## 2nd STEP: Moving Valve + Moving Mesh



volume fractions

0 0.2 0.4 0.6 0.8 1



absolute velocity (m/s)

0. 0.05 0.1 0.15 0.2



## PRO

- suitable to model the impact of the valve
- straight forward implementation
- robust and fast

## CONTRA

- not physical in the direct vicinity of the valve



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