# A new approach for Kirchhoff-Love plates and shells

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Kirchhoff-Love plates and shells

Joint work with Wolfgang Krendl, Katharina Rafetseder

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Extension to Mindlin-Reissner plates?

3 Extension to Kirchhoff-Love shells

4 Numerical results





Figure: Geometry of a shell

Unknown: covariant components of the displacement

 $\boldsymbol{u} = (u_i)$ 

plate: *S* is flat, for simplicity:  $\theta = Id$ pure bending:  $u_3 = w$ .

dimensional reduction of 3D elasticity for  $\Omega = \omega \times (-t/2, t/2)$ 

find  $w \in W \subset H^2(\omega)$  which minimizes

$$J(w) = \frac{1}{2} \int_{\omega} \mathcal{C} \nabla^2 w : \nabla^2 w \, dx - \int_{\omega} g \, w \, dx$$

with

$$C\varepsilon = 2\mu \left(\varepsilon + \frac{\nu}{1-\nu} (\operatorname{tr} \varepsilon) \boldsymbol{I}\right)$$

(kinematic) boundary conditions:

 $\partial \omega = \gamma_{\mathbf{c}} \cup \gamma_{\mathbf{s}} \cup \gamma_{\mathbf{f}}$ 

if the plate is simply supported on  $\gamma_c$ free on  $\gamma_f$ 

 $W = \{ v \in H^2(\omega) : v = 0, \partial_n v = 0 \text{ on } \gamma_c, v = 0 \text{ on } \gamma_s \}.$ 

variational formulation: find  $w \in W$  such that

$$\int_\omega \mathcal{C} 
abla^2 w : 
abla^2 v \ dy = \int_\omega g \ v \ dx \quad ext{for all } v \in W,$$

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quantities of interest: w vertical deflection  $M = -C\nabla^2 w$  bending moments

Mixed variational formulation: Find  $M \in V$  and  $w \in Q$  such that

$$\int_{\omega} \mathcal{C}^{-1} \mathbf{M} : \mathbf{L} \, dy + \int_{\omega} \mathbf{L} : \nabla^2 w \, dy = 0 \qquad \text{for all } \mathbf{L} \in \mathbf{V},$$
$$\int_{\omega} \mathbf{M} : \nabla^2 v \, dy \qquad = -\int_{\omega} g \, v \, dx \quad \text{for all } \mathbf{v} \in \mathbf{Q},$$

e.g., with the function spaces

$$V = L^2(\Omega)_{\text{sym}}, \quad Q = W.$$

An different choice for *Q* with less smoothness:

if

$$W = \{ v \in H^2(\omega) : v = 0, \partial_n v = 0 \text{ on } \gamma_c, v = 0 \text{ on } \gamma_s \},\$$

then

$$Q = \{ v \in H^1(\omega) : v = 0, \text{ on } \gamma_c, \quad v = 0 \text{ on } \gamma_s \} = H^1_{0, \gamma_c \cup \gamma_s}(\omega).$$

The choice for V: transposition method

$$b(\mathbf{v},\mathbf{L}) = \int_{\omega} \mathbf{L} : \nabla^2 \mathbf{v} \, d\mathbf{y} \equiv \langle B\mathbf{v},\mathbf{L} \rangle \longrightarrow \langle B^*\mathbf{L},\mathbf{v} \rangle,$$

where  $\langle B^* L, v \rangle$  is well-defined for  $v \in Q$ .

The linear operator **B**:

 $B: D(B) \subset Q \longrightarrow Y^*,$ 

here with

domain 
$$D(B) = W$$
 and  $Y = L^2(\Omega)_{sym}$ 

is a densely defined linear operator and unbounded operator in Q.

Its adjoint operator

$$B^*\colon D(B^*)\subset \mathbf{Y}\longrightarrow Q^*$$

is given by the identity

$$\langle B^* L, v \rangle = \langle Bv, L \rangle$$
 for all  $v \in W$ 

and all *L* from the domain

 $D(B^*) = \{ \boldsymbol{L} \in \boldsymbol{Y} \colon |\langle \boldsymbol{B} \boldsymbol{v}, \boldsymbol{L} \rangle| \leq c \, \|\boldsymbol{v}\|_Q \text{ for all } \boldsymbol{v} \in D(B) \}$ 

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standard integration by parts:

$$\int_{\omega} \mathbf{L} : \nabla^{2} \mathbf{v} \, d\mathbf{y}$$

$$= -\int_{\omega} \operatorname{Div} \mathbf{L} \cdot \nabla \mathbf{v} \, d\mathbf{y} + \int_{\partial \omega} (\mathbf{L}n) \cdot \nabla \mathbf{v} \, d\mathbf{s}$$

$$= -\int_{\omega} \operatorname{Div} \mathbf{L} \cdot \nabla \mathbf{v} \, d\mathbf{y} + \int_{\partial \omega} \mathbf{L}_{nt} \, \partial_{t} \mathbf{v} \, d\mathbf{s} + \int_{\partial \omega} \mathbf{L}_{nn} \, \partial_{n} \mathbf{v} \, d\mathbf{s}$$

$$? \qquad \mathbf{L}_{nn} = \mathbf{0} \text{ on } \gamma_{s} \cup \gamma_{f}$$

$$(\mathbf{L}_{nn} = \mathbf{L}n \cdot n, \quad \mathbf{L}_{nt} = \mathbf{L}n \cdot t)$$

 $\partial \omega = \gamma_s$ :

$$D(B^*) = \{ \boldsymbol{L} \in \boldsymbol{L}^2(\Omega)_{\text{sym}} : \text{ div Div } \boldsymbol{L} \in H^{-1}(\omega) \}$$

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#### Lemma

Let  $\mathbf{L} \in \mathbf{L}^2(\Omega)_{sym} \cap \mathbf{C}^1(\overline{\omega})$ . Then  $\mathbf{L} \in D(\mathbf{B}^*)$  if and only if

 $L_{nn} = 0$  on  $\gamma_s \cup \gamma_f$ 

#### and

 $\llbracket \boldsymbol{L}_{nt} \rrbracket_{x} = 0$  on interior corner points x of  $\gamma_{f}$ .

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AANMPDE 10 11 / 30

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Notations for  $B = \nabla^2 = \text{Grad grad}$ :

 $B^* = \operatorname{div}\operatorname{Div}, \quad D(B^*) = H(\operatorname{div}\operatorname{Div}, \omega; Q^*)_{\operatorname{sym}}.$ 

with

$$H(\operatorname{div}\operatorname{Div},\omega; Q^*)_{\operatorname{sym}} = \{ \boldsymbol{L} \in \boldsymbol{L}^2(\Omega)_{\operatorname{sym}} \colon \left| \int_{\omega} \boldsymbol{L} : \nabla^2 \boldsymbol{v} \, d\boldsymbol{y} \right| \le \boldsymbol{c} \, \|\boldsymbol{v}\|_{H^1(\omega)} \text{ for all } \boldsymbol{v} \in \boldsymbol{W} \}$$

Observe

$$H^1_0(\omega)_{
m sym} \subset H({
m div}\,{
m Div},\omega;\,Q^*)_{
m sym} \subset L^2(\Omega)_{
m sym}$$

Bernardi/Girault/Maday (1992), Z. (2015), Pechstein/Schöberl (2011), Rafetseder/Z. (2017)

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AANMPDE 10 12 / 30

Mixed variational formulation of

Find  $M \in V$  and  $w \in Q$  such that

$$\int_{\omega} \mathcal{C}^{-1} \mathbf{M} : \mathbf{L} \, d\mathbf{y} + \langle \operatorname{div} \operatorname{Div} \mathbf{L}, \mathbf{w} \rangle = 0 \qquad \text{for all } \mathbf{L} \in \mathbf{V},$$
$$\langle \operatorname{div} \operatorname{Div} \mathbf{M}, \mathbf{v} \rangle = -\int_{\omega} g \, \mathbf{v} \, d\mathbf{x} \quad \text{for all } \mathbf{v} \in \mathbf{Q},$$

with the function spaces

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$$V = H(\text{div Div}, \omega; Q^*)_{\text{sym}}, \quad Q = H^1_{0, \gamma_c \cup \gamma_s}(\omega),$$

equipped with norms  $\|v\|_Q = \|v\|_1$  and  $\|L\|_V = \|L\|_{\text{div Div};Q^*}$ , given by

$$\|\boldsymbol{L}\|^2_{\operatorname{div}\operatorname{Div};\boldsymbol{Q}^*} = \|\boldsymbol{L}\|^2_{\boldsymbol{L}^2(\Omega)} + \|\operatorname{div}\operatorname{Div}\boldsymbol{L}\|^2_{\boldsymbol{Q}^*}.$$

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#### Theorem

The mixed problem is well-posed in

$$\boldsymbol{V} \times \boldsymbol{Q} = \boldsymbol{H}(\operatorname{div}\operatorname{Div},\omega;\boldsymbol{Q}^*)_{\operatorname{sym}} \times \boldsymbol{H}^1_{0,\gamma_c\cup\gamma_s}(\omega).$$

#### Corollary

The primal variational problem and the mixed problem are equivalent:

- If w ∈ W solves the primal problem, then M = -C∇<sup>2</sup>w ∈ V and (M, w) solves the mixed problem.
- If (M, w) ∈ V × Q solves the mixed problem, then w ∈ W and w solves the primal problem.

#### Theorem

Let  $\omega$  be simply connected.  $M \in H(\text{div Div}, \omega; Q^*)_{\text{sym}}$  if and only if

 $\mathbf{M} = \mathbf{p} \mathbf{I} + \operatorname{sym}\operatorname{Curl} \phi$ 

with

$$p \in Q = H^1_{0,\gamma_c\cup\gamma_s}(\omega)$$
 and  $\phi \in \Psi_p$ 

with

$$\Psi_{p} = \left\{ \psi \in (H^{1}(\omega))^{2} : \langle \partial_{t}\psi, \nabla v \rangle_{\Gamma} = -\int_{\Gamma} p \, \partial_{n}v \, ds \quad \text{for all } v \in W \right\}$$

$$\phi = \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix}, \quad \text{sym}\,\text{Curl}\,\phi = \begin{pmatrix} \partial_2\phi_1 & \frac{1}{2}(\partial_2\phi_2 - \partial_1\phi_1) \\ \frac{1}{2}(\partial_2\phi_2 - \partial_1\phi_1) & -\partial_1\phi_2 \end{pmatrix}$$

Krendl/Rafetseder/Z. (2014,2016), Rafetseder/Z. (2017)

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Equivalent variational formulation in

 $p \in Q = H^1_{0,\gamma_c\cup\gamma_s}(\omega), \quad \phi \in \Psi_p \subset \left(H^1(\omega)\right)^2, \quad w \in Q = H^1_{0,\gamma_c\cup\gamma_s}(\omega).$ 

• Find  $p \in Q$  such that

$$\int_{\omega} 
abla p \cdot 
abla v \, dy = \int_{\omega} g \, v \, dx \quad ext{for all } v \in Q.$$

So For given  $p \in Q$ , find  $\phi \in \Psi_p$  such that

 $\int_{\omega} \mathcal{C}^{-1} \operatorname{sym}\operatorname{Curl} \psi : \operatorname{sym}\operatorname{Curl} \psi \, dy = \langle G[p], \psi \rangle \quad \text{for all } \psi \in \Psi_0.$ 

So For given  $M = pI + symCurl \phi$ , find  $w \in Q$  such that

$$\int_{\mathbb{R}^d} \nabla w \cdot \nabla q \, dy = \langle H[\textbf{\textit{M}}], q \rangle \quad \text{for all } q \in Q.$$

## Decomposition of the problem

Solution
for  $\phi^{\perp} = (-\phi_2, \phi_1)^T$  we obtain  $\int_{\omega} \hat{\mathcal{C}}^{-1} \varepsilon(\phi^{\perp}) : \varepsilon(\psi) \, dy = \langle \hat{G}[p], \psi \rangle \quad \text{for all } \psi \in \Psi_0^{\perp}$ 

with an appropriately rotated material tensor  $\hat{\mathcal{C}}^{-1}$  and

$$\varepsilon_{\alpha\beta}(\psi) = \frac{1}{2}(\partial_{\beta}\psi_{\alpha} + \partial_{\alpha}\psi_{\beta})$$

decomposition of the Kirchhoff model into

- Poisson problem
- Plane linear elasticity problem
- Poisson problem

Mindlin-Reissner model

$$J = \frac{1}{2} \int_{\omega} \left\{ \mathcal{C}\varepsilon(\theta) : \varepsilon(\theta) + \mu t^{-2} \|\nabla w - \theta\|^2 \right\} dx - \int_{\omega} g w dx \longrightarrow \min$$

Kirchhoff model:  $\theta = \nabla w$ .

$$J = \frac{1}{2} \int_{\omega} \left\{ \mathcal{C} \nabla^2 w : \nabla^2 w \right\} \, dx - \int_{\omega} g \, w \, dx \, \longrightarrow \min$$

decomposition of the Mindlin-Reisser model into

- Poisson problem
- Stokes-like problem
- Poisson problem

#### Brezzi/Fortin 1986

### Kirchhoff-Love shells

variational formulation of the linearized model:

find  $u \in W$  such that

$$\int_{\omega} \left[ t \, \mathcal{C} \varepsilon(\boldsymbol{u}) : \varepsilon(\boldsymbol{v}) + \frac{t^3}{12} \, \mathcal{C} \kappa(\boldsymbol{u}) : \kappa(\boldsymbol{v}) \right] \sqrt{a} \, d\boldsymbol{y} = \langle \boldsymbol{F}, \boldsymbol{v} \rangle \quad \text{for all } \boldsymbol{v} \in \boldsymbol{W}$$

with

$$\boldsymbol{W} \subset H^1(\omega) imes H^1(\omega) imes \left| H^2(\omega) \right|,$$

the membrane strain

$$oldsymbol{arepsilon}_{lphaeta}(oldsymbol{u}) = rac{1}{2}(oldsymbol{u}_{lpha|eta}+oldsymbol{u}_{eta|lpha}) - oldsymbol{b}_{lphaeta}\,oldsymbol{u}_3,$$

the bending strain

$$egin{aligned} \kappa_{lphaeta}(oldsymbol{u}) &= oldsymbol{u}_{3|lphaeta} - oldsymbol{b}^{\sigma}_{lpha}oldsymbol{b}_{\sigmaeta}oldsymbol{u}_3 + oldsymbol{b}^{\sigma}_{lpha}oldsymbol{u}_{\sigma|eta} + oldsymbol{b}^{\tau}_{eta}oldsymbol{u}_{ au|lpha} + oldsymbol{b}^{\tau}_{eta|lpha}oldsymbol{u}_{ au} \ &= oldsymbol{eta}_{lphaeta}oldsymbol{u}_3 + \dots . \end{aligned}$$

#### Kirchhoff-Love shells

quantities of interest:  $\mathbf{u}$  displacement  $\mathbf{M} = \frac{t^3}{12} C \kappa(\mathbf{u}) \sqrt{a}$  bending moments  $= \hat{C} \kappa(\mathbf{u})$ 

Mixed variational formulation: Find  $M \in V$  and  $u \in Q$  such that

$$\int_{\omega} \hat{\mathcal{C}}^{-1} \boldsymbol{M} : \boldsymbol{L} \, d\boldsymbol{y} \qquad -\langle \operatorname{div} \operatorname{Div} \boldsymbol{L}, \boldsymbol{u}_{3} \rangle - \int_{\omega} \boldsymbol{L} : \boldsymbol{\kappa}^{1}(\boldsymbol{u}) \, d\boldsymbol{y} = 0$$
$$-\langle \operatorname{div} \operatorname{Div} \boldsymbol{M}, \boldsymbol{v}_{3} \rangle - \int_{\omega} \boldsymbol{M} : \boldsymbol{\kappa}^{1}(\boldsymbol{v}) \, d\boldsymbol{y} - \int_{\omega} t \, \mathcal{C} \boldsymbol{\epsilon}(\boldsymbol{u}) : \boldsymbol{\epsilon}(\boldsymbol{v}) \, \sqrt{a} \, d\boldsymbol{y} \qquad = -\langle \boldsymbol{F}, \boldsymbol{v} \rangle$$

for all  $L \in V$  and  $v \in Q$ , with the function spaces

 $V = H(\text{div Div}, \omega; Q_3^*)_{\text{sym}},$  $Q = Q_1 \times Q_2 \times Q_3,$ 

where  $Q_i \subset H^1(\omega)$ .

benchmark examples of the shell obstacle course:

- Scordelis-Lo roof
- pinched cylinder
- pinched hemisphere

midsurface represented as B-spline surfaces

approximation spaces:

- equal-order approximation for  $u_i$ , p,  $\phi$
- single patch: B-splines of maximum smoothness
- multipatch: C<sup>0</sup>

implementation in G+Smo (C++ library for IgA)

Example: Scordelis-Lo roof



Figure: Geometry and load

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AANMPDE 10 22 / 30

#### Example: Scordelis-Lo roof



#### Figure: Results, 1 patch

#### Example: Scordelis-Lo roof



#### Figure: Results, 4 patches

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AANMPDE 10 24 / 30

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Example: pinched Cylinder



Figure: Geometry and load

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AANMPDE 10 25 / 30

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Example: pinched Cylinder



#### Figure: Results, 4 patches

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AANMPDE 10 26 / 30

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Example: pinched Hemisphere



Figure: Geometry and load

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AANMPDE 10 27 / 30

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#### Example: pinched Hemisphere



#### Figure: Results, 4 patches

Kirchhoff-Love plates and shells

AANMPDE 10 28 / 30

- Kirchhoff plate bending problems and similar 4-th order problems can be decomposed in three (consecutively to solve) second-order problems.
- extension to Kirchhoff-Love shell: formulation in  $H^1$  spaces.
- work in progress:
  - extend the numerical analysis from Kirchhoff plates to Kirchhoff-Love shells.
  - behavior of equal-order discretization method w.r.t. membrane locking
  - extension to Mindlin-Reissner plates and shells

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A decomposition result for biharmonic problems and the Hellan-Herrmann-Johnson method.

ETNA, Electron. Trans. Numer. Anal., 45:257–282, 2016.

#### K. Rafetseder and W. Z.

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